# The Universal Speed Limit as an Attribute of the Space-Time Symmetry

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# Abstract-Relativistic kinematic effects are deduced immediately from the space-time symmetry.

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The theory of relativity [1, 2], that began primarily with explanations of concrete physical effects, is now assumed [3, 4] to be just a reflection of the space-time symmetry resulting in the relativistic coordinate-time transform [1, 2] without using a priori limitation on speeds of material bodies. Historically, the latter approach was initiated by Proclus [5], who abolished the fifth element of Aristotle (the ether) and, so, converted both the motion and the rest from absolute to relative ones.

The present comment to [3-5] reproduces the author's talk [6] and exercises with physics students. Following the antique methodology [5], start with definitions. Launch two identical crystals into the infinite vacuum along a common direct line (Fig. 1). In each reference frame, measure any distance,  $\xi$  or  $\xi'$ , with an unit equal to N crystal periods, and measure any time interval,  $\tau$  or  $\tau'$ , with an unit equal to M electron "rotations" around any local nucleus.



Fig. 1 Inertial reference frames representing crystals moving with different speeds (denoted with different thick arrows). From the external viewpoint, the identical crystals seem squeezed to different lengths - because of different Lorenz "length contractions" [1, 2]

To meet the aesthetic "ancient-Greek" symmetry ( $\sigma \nu \mu \epsilon \tau \rho \epsilon i \nu =$  symmetrein means "to measure together"), assume the ratio *N/M* being adjusted so that both coordinates  $\xi, \xi'$  and times  $\tau, \tau'$  are measured with a common uni-dimensional etalon<sup>\*</sup>. Under this convention, any registration of any event (indicated by a flash in Fig. 1 and Fig. 2) in any of the reference frames should be invariant relative to the commutation

$$\begin{split} \xi \to \tau & \xi' \to \tau' \\ \tau \to \xi & \tau' \to \xi' \end{split} \tag{1}$$

In addition, with account of the counter-directional mutual reciprocity of the identical reference frames, the space-time isotropic interrelation is to be invariant relative to the commutation

$$\begin{split} \xi \to -\xi' & \xi' \to -\xi \\ \tau \to \tau' & \tau' \to \tau \end{split} \tag{2}$$

Put  $\xi_0 = \tau_0 = \xi_0' = \tau_0' = 0$  for a primary event and, to meet the symmetry invariance conditions (1)-(2), interrelate subsequent events  $(\xi, \tau \quad \xi', \tau')$  by equations

$$\begin{aligned} \xi &= \gamma \xi' + \kappa \tau' \quad \Longrightarrow \quad \xi' = \gamma \xi - \kappa \tau \\ \tau &= \kappa \xi' + \gamma \tau' \quad \longleftrightarrow \quad \tau' = -\kappa \xi + \gamma \tau \end{aligned}$$
(3)

<sup>\*</sup>At this point, progressive modern students propose: "Let both N periods and M rotations be equal to 1\$".

By putting the right equations into the left ones and by putting the left equations into the right ones, derive equations of the following type:  $\xi = \gamma(\gamma\xi - \kappa\tau) + \kappa(-\kappa\xi + \gamma\tau) = (\gamma^2 - \kappa^2)\xi$ , which result in the common condition

$$\gamma^2 - \kappa^2 = 1 \tag{4}$$

(converted to the Pythagoras theorem in Appendix 1).

By putting  $\xi' = 0$  or  $\xi = 0$  into Eqs. (3), find the relative velocity of the reference frames

$$\beta = \pm \kappa / \gamma. \tag{5}$$

By using Eq. (5), convert the cyclic recurrence condition (4) to the form

$$\gamma = \left(1 - \beta^2\right)^{-1/2},\tag{6}$$

resulting in the limitation

$$|\boldsymbol{\beta}| \le 1 \tag{7}$$

for relative velocities of reference frames and, so, of any material bodies. Thus, the space-time-symmetry postulate (1)-(2) has given additional proof (7) for the 12th theorem of Proclus [5]: "Εν πεπερασμένω χρόνω το άπειρον κινείσθαι ουκ έστιν" = "During a limited time it is not possible to go an infinite distance".

The space-time transform formulas (3)-(4) are followed by popular effects of the relativistic kinematics [1, 2]:

- substitutions  $\tau=0$  or  $\tau'=0$  into Eqs. (3) give, correspondingly,  $\xi = \xi'/\gamma$  or  $\xi' = \xi/\gamma$ , which is called the Lorenz "length contraction" (shown in Fig. 1);
- substitutions  $\xi'=0$  or  $\xi=0$  into Eqs. (3) give, correspondingly,  $\tau'=\tau/\gamma$  or  $\tau=\tau'/\gamma$ , which is called the "twin paradox";
- if a body is moving with a velocity β'<sub>b</sub> = ξ'/τ' relative to the {ξ', τ'} reference frame, then according to Eqs. (3)
  the body velocity relative to the {ξ, τ} frame is

$$\beta_b = \xi / \tau = \frac{\beta_b' + \beta}{1 + \beta \beta_b'},\tag{8}$$

which is called the "relativistic velocity summation" - consistent with the Proclus speed limitation (7);

• according to limitation (7) and Eq. (8), in all inertial reference frames all small vacuum perturbations propagate with a common – ultimate – speed equal to 1 (exemplified with the electromagnetic pulse propagation [1, 7] - Appendix 2).

In Appendix 3, formulas (3)-(4) are converted to the conventional relativistic kinematic transform [1, 2]. However, note that, economically compared to [1, 2], the commutation Eqs. (3)-(4) have been above deduced from the only postulate - of the aesthetic space-time-symmetry [3-5].

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### APPENDIX 1 COORDINATE-TIME TRANSFORM IN COMPLEX VARIABLES

It may be pedagogical to undertake the Minkowski's [1, 2] change of variables  $\hat{\tau} = i\tau$ ,  $\hat{\tau}' = i\tau'$ ,  $\hat{\kappa} = i\kappa$  converting the Eqs. (3) to the form

$$\begin{aligned} \xi &= \gamma \xi' - \hat{\kappa} \hat{\tau}' \\ \hat{\tau} &= \hat{\kappa} \xi' + \gamma \hat{\tau}' \end{aligned} \iff \begin{aligned} \xi' &= \gamma \xi + \hat{\kappa} \hat{\tau} \\ \hat{\tau}' &= -\hat{\kappa} \xi + \gamma \hat{\tau} \end{aligned} \tag{A1.1}$$

and the cyclic recurrence condition (4) to the form

$$\det \begin{vmatrix} \gamma & \mp \hat{\kappa} \\ \pm \hat{\kappa} & \gamma \end{vmatrix} = 1.$$
(A1.2)

These formulas may be illustrated with Fig. 2: taking into account the Euclidean similarity of the triangles, the "Descartes coordinates"  $\xi$ ,  $\hat{\tau}$  and  $\zeta'$ ,  $\tau'$  are mutually turned at the angle  $\alpha = \arcsin \hat{\kappa} = \arccos \gamma$ . Correspondingly, the Eq. (A1.2) takes the form of "Pythagoras theorem"  $\cos^2 \alpha + \sin^2 \alpha = 1$ , where, the angle  $\alpha$  being imaginary, the "cathetuses"  $\cos \beta = \gamma$  and  $\sin \alpha = \hat{\kappa}$  may be longer than the "hypotenuse" 1.



Fig. 2 Rotation of Descartes coordinates  $\xi, \hat{\tau}$  and  $\xi'$  ,  $\hat{\tau}'$  .

# APPENDIX 2 VACUUM PERTURBATIONS

Take the wave equations [1, 7]

$$\frac{\partial^2 \varphi}{\partial \xi^2} = \frac{\partial^2 \varphi}{\partial \tau^2} \Leftrightarrow \frac{\partial^2 \varphi}{\partial {\xi'}^2} = \frac{\partial^2 \varphi}{\partial {\tau'}^2}$$
(A2.1)

describing an electromagnetic perturbation  $\varphi$  of the vacuum (the unitary coordinate and time in these equations are related to the "usual" ones in Appendix 3). These equations satisfy the double symmetry conditions (1)-(2) and, so, should be invariant relative to the space-time commutation Eqs. (3)-(4). For checking, use Eqs. (3) to derive

$$\frac{\partial^2 \varphi}{\partial \xi^2} = \gamma^2 \frac{\partial^2 \varphi}{\partial \xi'^2} - 2\gamma \kappa \frac{\partial^2 \varphi}{\partial \xi' \partial \tau'} + \kappa^2 \frac{\partial^2 \varphi}{\partial \tau'^2},$$

$$\frac{\partial^2 \varphi}{\partial \tau^2} = \gamma^2 \frac{\partial^2 \varphi}{\partial \tau'^2} - 2\gamma \kappa \frac{\partial^2 \varphi}{\partial \xi' \partial \tau'} + \kappa^2 \frac{\partial^2 \varphi}{\partial \xi'^2},$$
(A2.2)

put these derivatives into the left (A2.1) and see how the cyclic recurrence formula Eq. (4) converts the left (A2.1) to the right (A2.1).

## APPENDIX 3 INTERRELATION BETWEEN UNITARY AND "USUAL" SPACE-TIME SCALES

To relate the uni-dimensional symmetrized coordinate  $\xi$  and time  $\tau$  to their "usual" (measured, for instance, with meters and seconds) equivalents  $\chi$  and t, undertake the linear change of variables

$$x = a\xi, \quad t = b\tau \tag{A3.1}$$

converting the "ancient-Greek" commutation Eqs. (3)-(4) to the "usual" Lorenz transform [1, 2]

$$\begin{array}{l}
x = \gamma(x' + vt') \\
t = \gamma(t' + vx'/c^2) & \Leftrightarrow \begin{array}{l}
x' = \gamma(x - vt) \\
t' = \gamma(t - vx/c^2) \\
\gamma = 1/\sqrt{1 - v^2/c^2}
\end{array}$$
(A3.2)

where  $v = \beta c$  is the "usual" relative velocity of reference frames and a/b = c is the "usual" speed limit common for all material bodies. Measured at the end of the previous millennium, this absolute limit - the speed of light – was proved to be  $3 \cdot 10^8 m / s$  [1, 2].

### REFERENCES

- [1] W. Pauli, *Theory of Relativity*, Dover Publications, 1981.
- [2] R. Feynman, R. Leighton, and M. Sands, The Feynman Lectures on Physics, 2006, ISBN 0-8053-9045-6.
- [3] N. Mermin, "Relativity Without Light," Am. J. Phys., vol. 52, pp. 119-124, 1984.
- [4] M. Feigenbaum, The Theory of Relativity Galileo's Child, arXiv: 0806.1234v1, 2008.
- [5] Πρόκλου Διαδόχου Αυκίου, "Στοιχείωσις Φυσική", Elements of Physics, A. Ritzenfeld (ed./transl.), Leipzig: Teubner, 1912.
- [6] M. Petelin, "A Compact Derivation of the Relativistic Time-Coordinate Transform," *Proceedings of IW-FIRT 2012*, Fukui, Japan, P. 1, 2012.

[7] W. Voigt, Ueber das Dopplersche Prinzip, Goett. Nachr. s 41, 1887.